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For the CDF
Collaboration

Minimum Bias Studies at CDF and Comparison with MonteCarlo

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This talk

- Set of measurements of inelastic non-diffractive particle production in $p\bar{p}$ interactions collected with a MB trigger:
 - Charged particle p_T differential cross section
 - Previous measurements (CDF 1988) extended in range and precision
 - Correlation of charged particle $\langle p_T \rangle$ with their multiplicity in the event
 - ΣE_T differential cross section
- Systematic comparison with MC at hadron level
- One example of possible outcome (top mass)
- Point out experimental issues and problems

Introduction

- MB data contains (in principle) all type of interactions proportionally to their natural production rate.
- At 2 TeV is a mixture of soft and hard interactions.
 - $\sim 1\%$ events with jet of $E_T > 10$ GeV, $\sim 25\%$ jet $E_T > 3$ GeV
- It is the most suitable sample to study at the same time:
 - Non-perturbative contributions
 - Low- Q^2 parton interactions
 - beam-beam remnants
 - MPI
 - Other mechanisms (correlations?)
 - Interplay of soft and hard
- Important in itself and for precision measures of hard scattering processes (subtraction of soft effects)

} U.E.

By comparison to future LHC measurements, we will gather information on the energy dependence of the MPI process

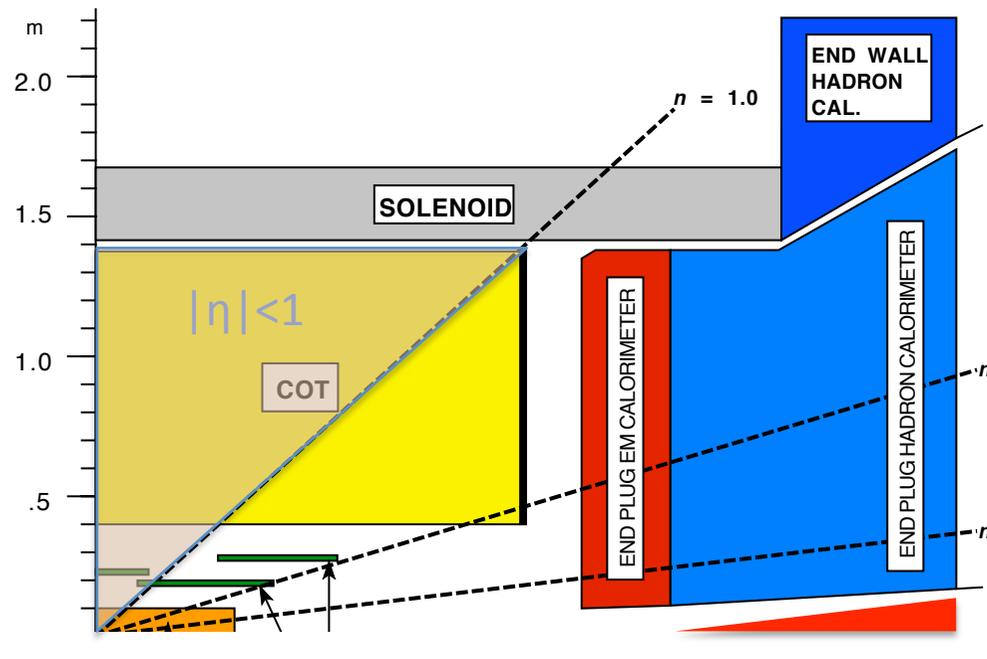
Especially at LHC because of huge pile-up !

MB and MC

- No generator correctly reproduces all the observables at the same time, but most may be tuned to be ok on each *single* one.
- In Pythia the description of MB rely on a \hat{p}_T cut-off that :
 - Regulates the 2-2 parton perturbative Xs at low momenta, but also...
 - ...the additional parton-parton scatterings (MPI)
 - Plus all the stuff we'll see in other presentations...
- Pythia tries to describe the largest possible set of MB distributions

A “small-bias” trigger

- “Minimum Bias” is defined by its trigger !

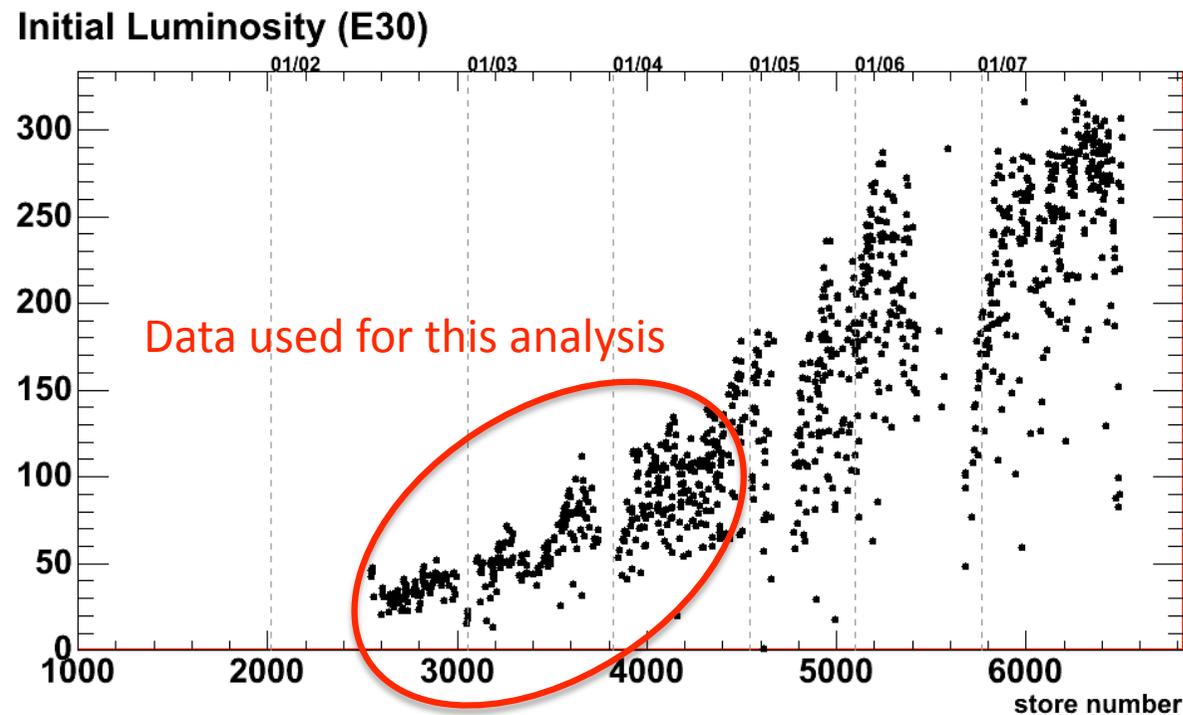


- CDF triggers MB with forward particles in $3.7 < |\eta| < 4.7$ ($< 3\text{deg}$)
- Requires coincidence of both sides (forw+backw)
- Implemented with Cerenkov counters (CLC)

- The trigger is biased so to favor high p_T interactions
- Will affect the *shape* of inclusive distributions
- Note that the observables studied are in the central region ($|\eta| < 1.0$)

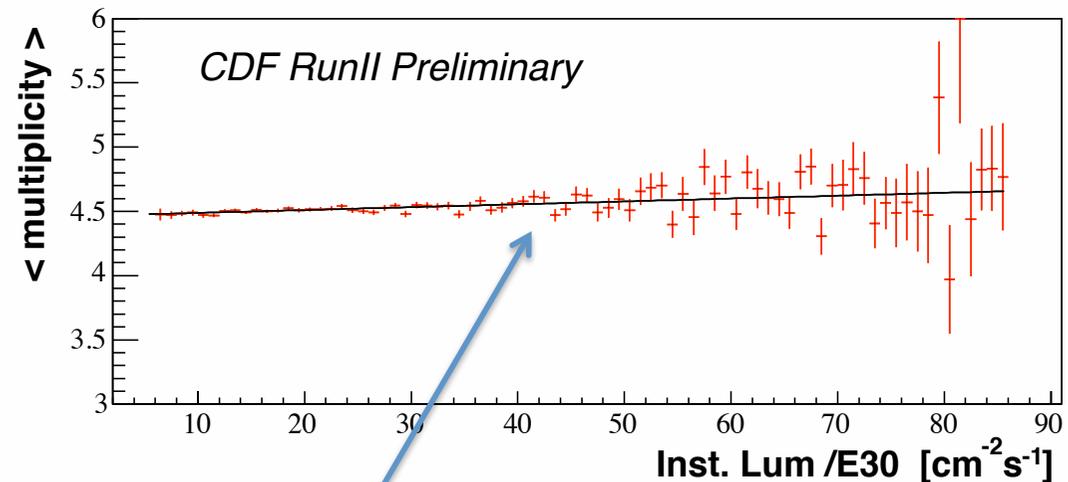
The data Sample

- 506 pb^{-1} (Tevatron runs during 2002 to 2004)
- Max inst. Lum: $90 \times 10^{30} \text{ cm}^{-2}\text{s}^{-1}$ (<50 for ΣE_T)
- Average inst. Lum: $20 \times 10^{30} \text{ cm}^{-2}\text{s}^{-1}$ (<17 for ΣE_T)



The Event Selection and Background

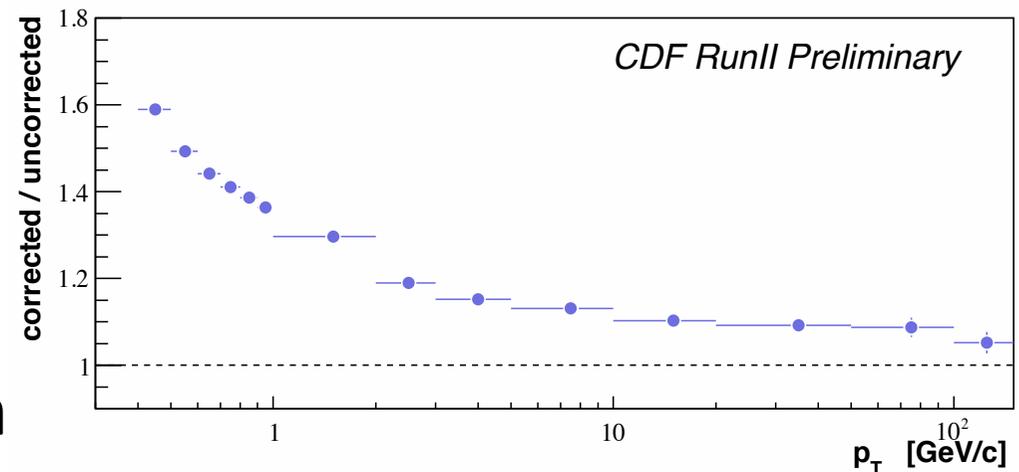
- Event cuts:
 - $|Z_{\text{vertex}}| < 40$ cm
 - Limited vertex asymmetries in η
 - No pile-up detected
- Vertex acceptance is function of the total activity in the crossing
- Backgrounds:
 - Pile-Up ($\sim 3\%$ of MB)
 - Diffractive interactions ($\sim 3.5\%$ of MB)



The average particle multiplicity grows with the Inst.Lum. because of undetected pile-up events. The resolution for separating primary vertices in Z is ~ 3 cm

Tracking efficiency

- ϵ is function of p_T , N_{ch} , # of jets...
 - In the same event the probability for a particle to be observed depends whether is inside a jet or not (and on the jet E_T)
 - Cannot analyze as function of all kinematic variables, must rely on MC for most
- Overall correction computed with MC:
 - ϵ is the largest contribution: ($\sim 70\%$ at $p_T=0.4$, $\sim 92\%$ at $p_T=5$ GeV/c)
 - includes correction for
 - Contamination of second.
 - Particle decays, conversions..
 - Trigger + event acceptance
- Full detector simulation and event reconstruction
- Accepted region: $p_T > 0.4$ GeV/c and $|\eta| < 1.0$



Single Particle p_T differential Xs

$$E \frac{d^3\sigma}{dp^3} = \frac{d^3\sigma}{p_T \Delta\phi \Delta y dp_T} = \frac{N_{ch} / (\epsilon \times A)}{L p_T \Delta\phi \Delta y dp_T}$$

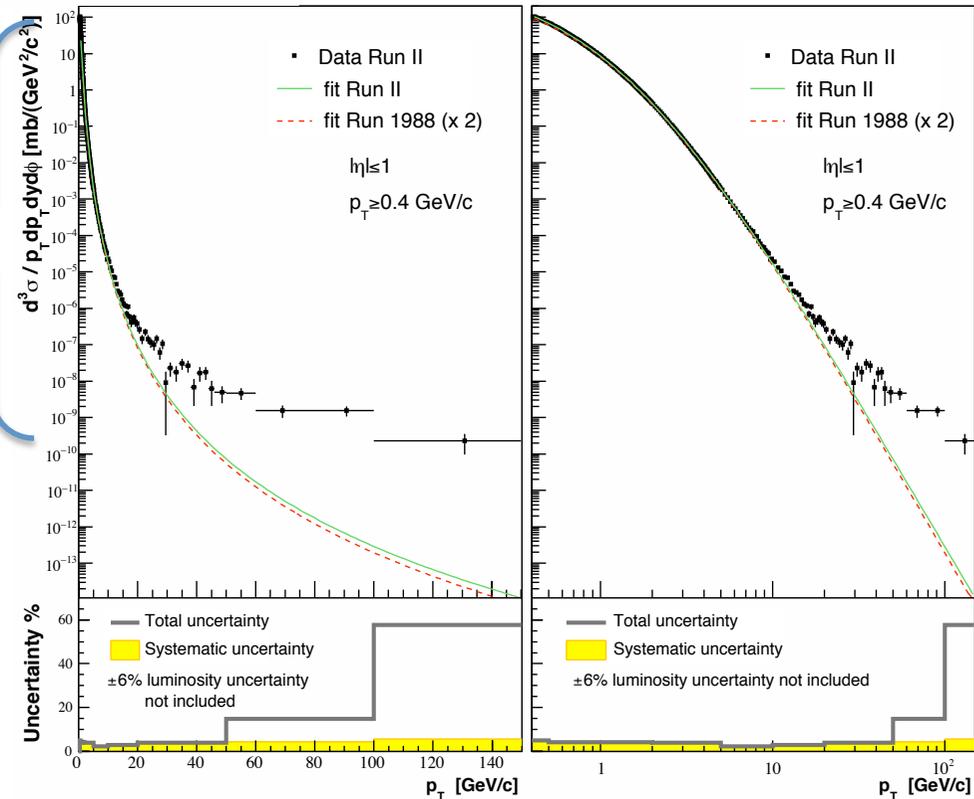
$A(Z_{vertex}, \text{PileUp})$

- Modeling of the distribution:

$$f = A \left(\frac{p_0}{p_T + p_0} \right)^n$$

$10^2 - 10^{-9}$

- High p_T tail orders of magnitude higher than fit extrapolation
- Same fit as with 1800 GeV
- New parameterization ?

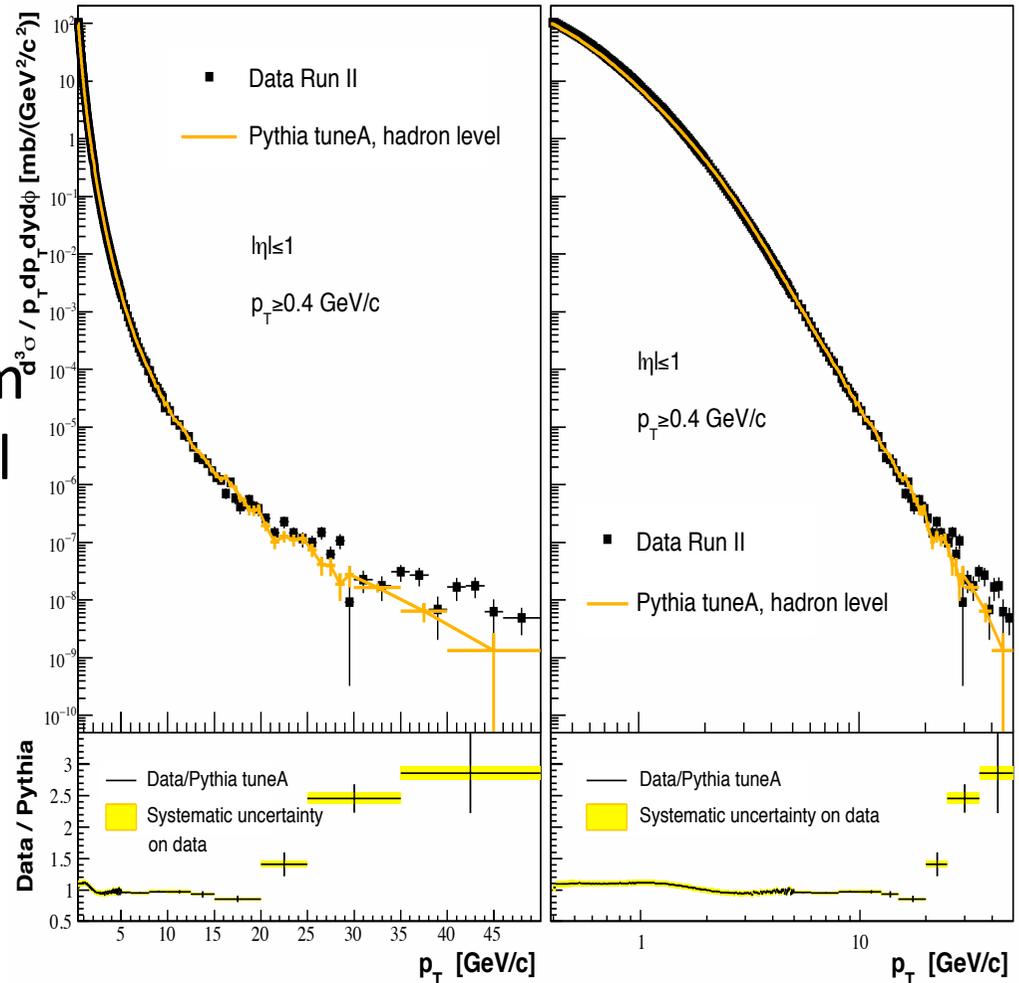
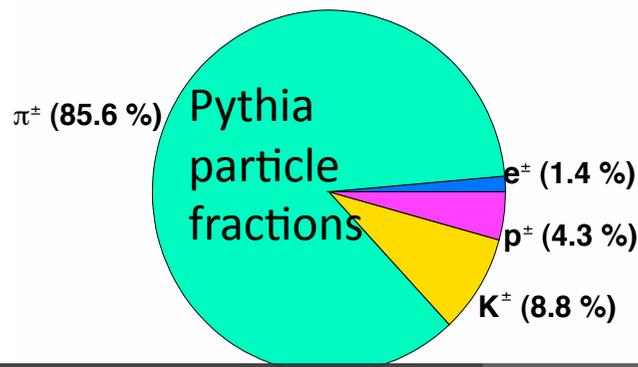


$X_s(2\text{TeV}) \sim 1.04 X_s(1.8\text{TeV})$

	p_0	n	p_T range (GeV/c)	χ^2/dof
Run 0, 1800 GeV	1.29 ± 0.02	8.26 ± 0.08	0.4 - 10.	102/64
Run 0, 1800 GeV	1.3 fixed	8.28 ± 0.02	0.4 - 10.	103/65
Run II, 1960 GeV	1.230 ± 0.004	8.13 ± 0.01	0.4 - 10.	352/192

P_T Xs, compare to Pythia

- Data has much more particle production in $p_T > 20$ GeV
- TuneA produces no particles at all above ~ 50 GeV/c
- Note that Δy is computed from $\Delta \eta$ assuming the π mass for all tracks
- Bias $\sim 5\%$ at $p_T = 0.4$ GeV/c dependent on the (unknown) fractions of particle species

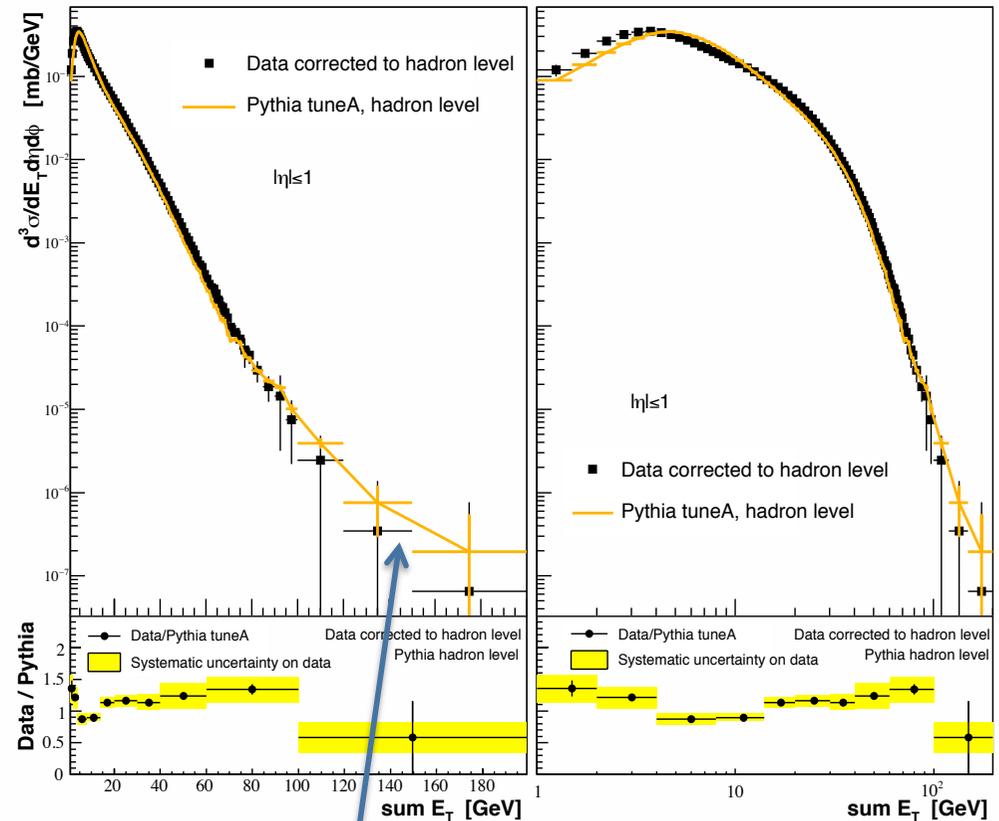


Overall 6% systematic from Luminosity measurement

Sum E_T differential Xs

$$\frac{d\sigma}{\Delta\phi \Delta\eta dE_T} = \frac{N_{ev}}{\text{Lum} \Delta\phi \Delta\eta dE_T}$$

- First attempt of really inclusive measure with neutral particles
- ΣE_T integrates all undetected pile-up: large systematic also at low lum ($\sim 3\%$)
- Different Pythia tunings (A,DW) give different cal responses: systematic ~ 5 to 15%
- + overall systematic 6% due to luminosity measurement

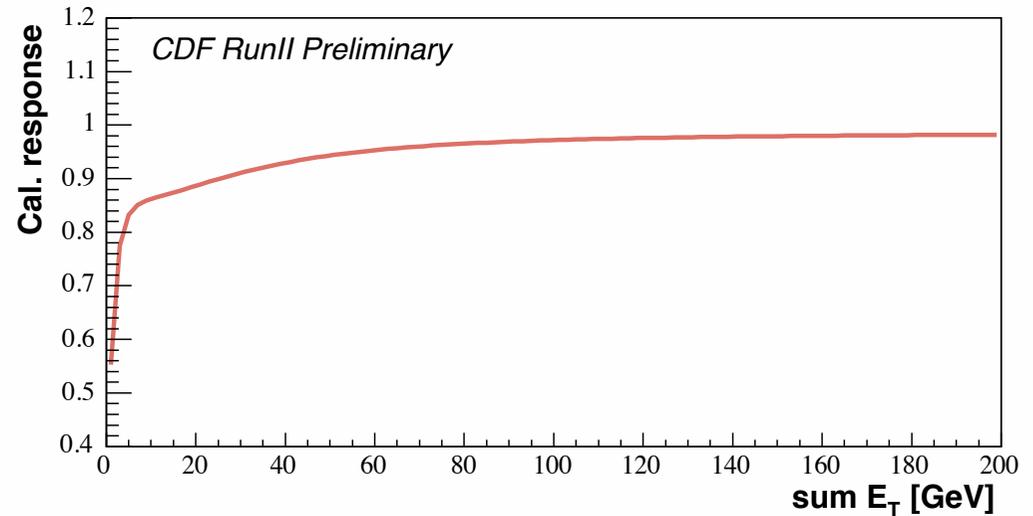


Excess of energy in the UE in high p_T jets wrt data ?

Corrections to Sum E_T Xs

■ Many step correction:

- Single tower relative cal response $f(\eta, Z_{\text{vertex}})$
- Absolute response to ΣE_T
- Acceptance vs Z_{vertex}
- Pile-Up
- Undetected low- p_T charged particles →
- Trigger and vertex acceptance
- Unfolding (spread of events due to finite energy resolution)



Cut at $p_T \sim 0.3$ GeV/c at CDF
1 to 2 GeV/c at CMS and Atlas

$$N_{ev}^{corrected} = N_{ev}^{raw} \left(\frac{\Sigma E_T}{C_{tower-\eta} C_{absolute} A_{vertex-Z}} + C_{low-p_T} \right) \frac{C_{pile-up} \times C_{unfolding}}{A_{trigger \text{ and } vertex}}$$

Charged particle $\langle p_T \rangle$ vs Multiplicity

$$\langle p_T \rangle (N_{ch}) = \frac{\sum_{ev} \sum_i^{N_{ch}} p_T^i}{N_{ev}^{N_{ch}} \times N_{ch}}$$

■ Sensitive to the various components of MB:

- Measure of the amount of hard vs soft
- Sensitive to the modeling of MPI, eg . 'Color Reconnections'
- Most poorly reproduced by MC
- Probably best current model is Pythia tuneA, but new tunings with ad-hoc CR + much more are being studied

P.Skands and D.Wicke, 2007:

"...models which successfully describe the $\langle p_T \rangle (N_{ch})$ distribution such as R.Field 'TuneA' ... do so by incorporating very strong ad hoc correlations between final state partons from different interactions...

..Simultaneous agreement with $\langle p_T \rangle (N_{ch})$ is only obtained from the models incorporating non-trivial color correlations..."

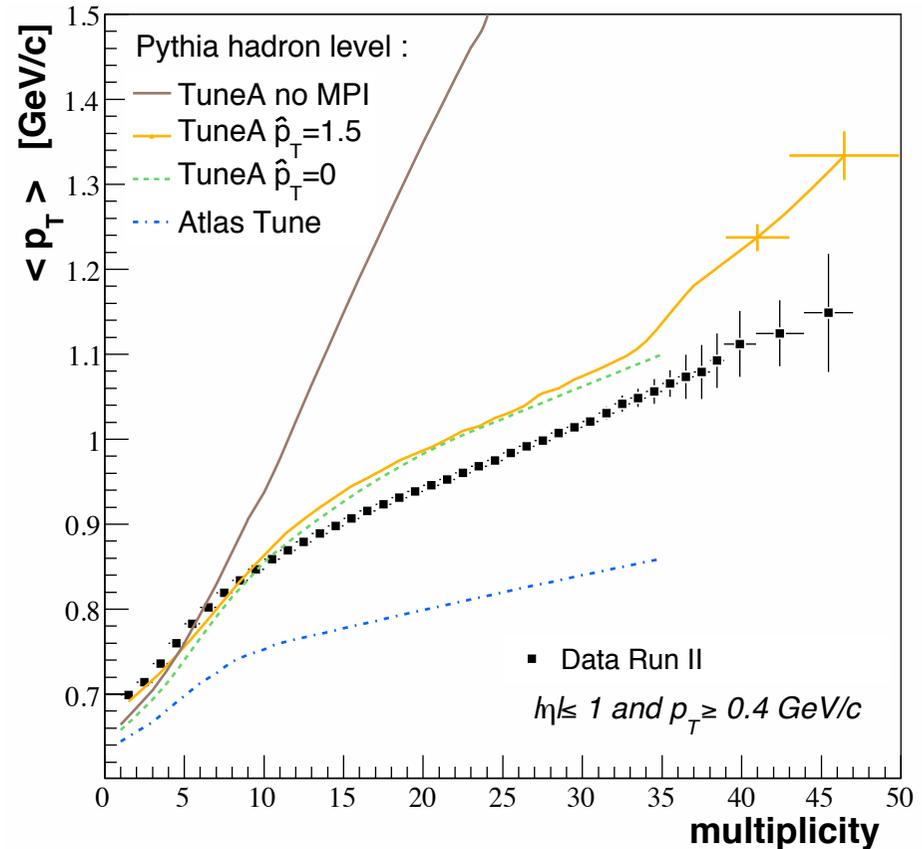
$\langle p_T \rangle (N_{ch})$

- Merged data at high N_{ch} from a dedicated “high-multiplicity” trigger

- ToF at L1
- Track reconstruction at L3
- Double the stat in $N_{ch} > 24$
- Total uncertainty
~2% (6% for $N_{ch} > 40$)

- TuneA: fairly good
 - TuneA no MPI: too hard
 - Atlas tune: too soft
- (tunes by courtesy of R.Field)

- MPI mechanism necessary to reproduce $\langle p_T \rangle (N_{ch})$
- $\langle p_T \rangle (N_{ch})$ useful to tune MPI



Outcomes: an example

- New Pythia tuning based on these data

- New parton shower and ISR/FRS models
- Ad hoc 'color reconnection' model

- Effect on the had TOP mass measure

- Produce a sample of $t\bar{t} \rightarrow Wb+Wb \rightarrow 6$ jets
- Compare to previous tuneA
- M_{top} variation ~ 0.6 to 0.8 GeV

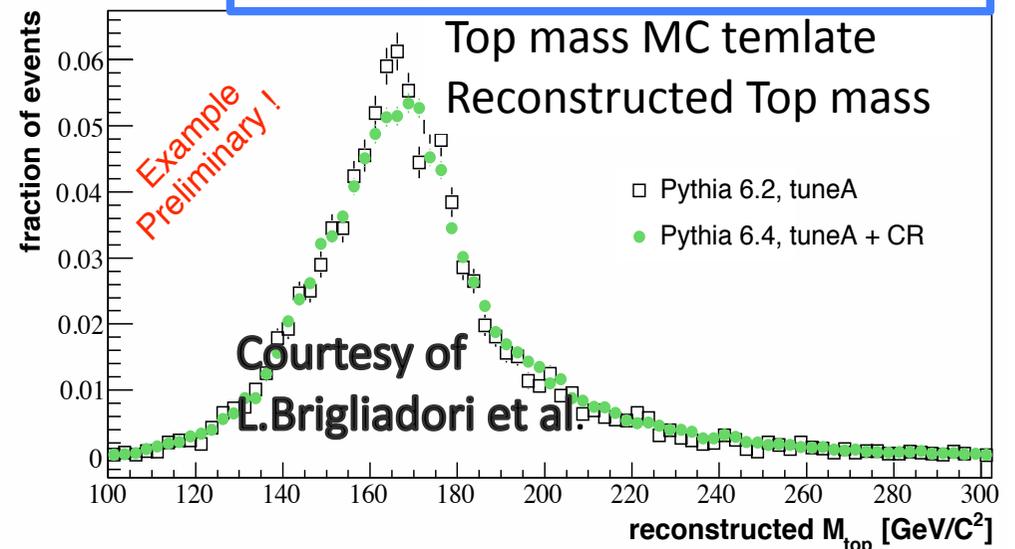
- But first we need to disentangle perturbative

effects (jet shapes should be sensible to parton shower) and check the effect on

- other top channels
- dijets
- JES...

All-hadronic TOP:

- Many templates produced with varying M_{top}
- Background is added
- Templates are parameterized
- Parameters studied vs M_{top} , compared to data best set chosen



Experimental issues: caveats for LHC

- By the experience gained at CDF, some effects affect the analyses of MB:
 - Undetected pile-up becomes systematic uncertainty
 - Need excellent separating resolution of PV (timing may be a good handle)
 - Watch for biases introduced by the selection of the PV (no algorithm can be totally unbiased)
 - MB triggers like CDF are actually biased to favour hard interactions (eg, Atlas may incur in the same problem)
 - Low- p_T particle reconstruction usually neglected (~ 1 GeV/c ?)
 - At LHC the high B fields will hide a large part of the spectrum !
 - Care needed in tracking inside/outside jet cones
 - Do not forget the rest of the world (neutrals) and P.ID.
- MB is (also) useful to the understanding of the detector and to calibrate tools of high- p_T analyses (eg. JES, ME_T ...)

Conclusions

- MB is a mixture of different processes:
 - It's difficult to simulate, but...
 - It's *the* sample where is possible to understand how all the processes involved overlap and interact
- We have provided a set of high precision measurements of the MB final state at 2 TeV
- Few immediate observations:
 - Pythia reproduces the particle p_T spectrum within 10% only at $p_T < 20$ GeV/c
 - $\langle p_T \rangle (N_{ch})$ strongly favors models with MPI and put constraints on the MPI model itself
- Hopefully these data will lead to improved models in time for the first LHC data

BACKUP SLIDES

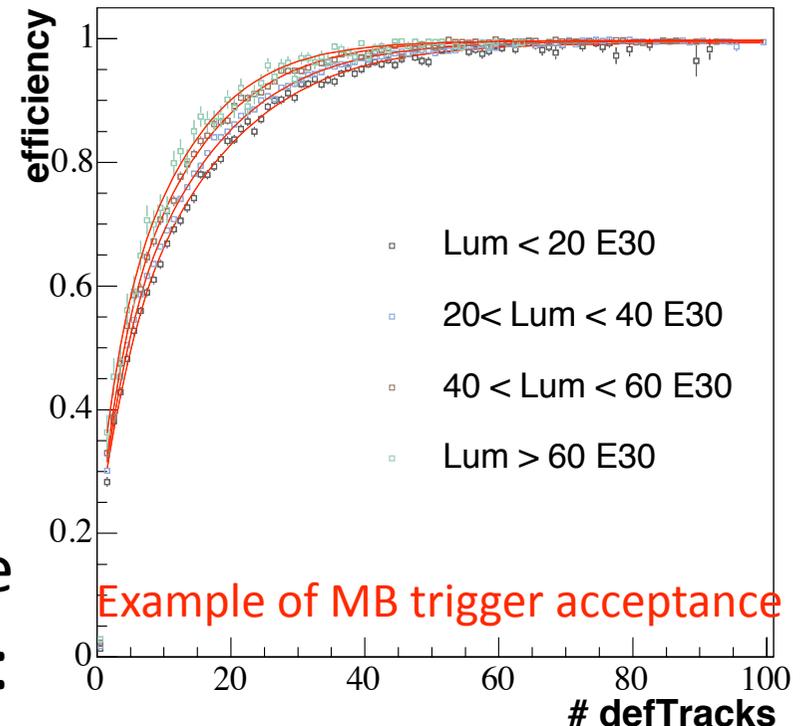
CDF Vertex & Trigger acceptance

- Trigger: acceptance is function of many variables

- # primary detected vertices, Luminosity, # tracks, CLC calibration...
- In short, acceptance increases with the total activity in the *crossing*. Plateau at ~98%
- Measurable with 0-Bias sample

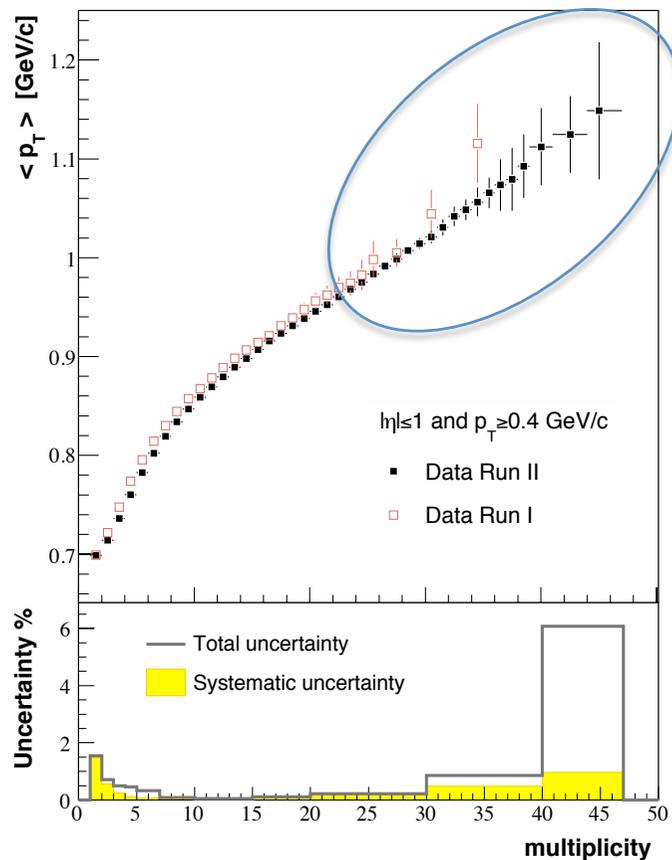
- Primary Vertex (after trigger):

- Clustered by the tracking system
- Flat in $|Z_{\text{vertex}}| < 40$ cm, outside affected by tracking inefficiencies
- Depends on #PV, # tracks, Luminosity



$\langle p_T \rangle (N_{ch})$ – comparison to Run I

- This measurement was published by CDF in 2002 (PRD D65 072005)



Correction procedure

1. Tracking efficiency: correct the $\langle p_T \rangle$ for each reconstructed N_{ch}
2. Smearing: generate matrix of $P(N_r, N_g)$
Then apply this unfolding factor bin by bin

$$\langle p_T \rangle_{n_r=m} = \sum_i^{n_g} (\langle p_T \rangle_{n_r=i} \times P_{n_r=m}^{n_g=i})$$

$\langle p_T \rangle$ at N_r

$\langle p_T \rangle$ at N_g

(.. works as long as $\langle p_T \rangle$ at $N_{ch} = \text{gen}$ is the same that at $N_{ch} = \text{rec}$)